

# The work input to a granular material

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## INTRODUCTION

Some recent theoretical models for soils are based on hypotheses about the rate at which the input work to the soil is either stored or dissipated. It is therefore necessary to evaluate the rate of work input to the soil in terms of the stresses, strain rates and other variables. For a single-phase material the power input per unit volume is simply the product of the stresses and the strain rates; but this result does not apply for the two-phase material, such as a saturated soil, where both the stresses within the two phases and the velocities of the two materials will be different. The total power input per unit volume to a soil must therefore be derived by considering the rate at which all the forces on both the soil grains and the pore fluid do work.

The principle of effective stress as described by Terzaghi (1943) states that the mechanical behaviour of a soil is governed by the difference between the total stress and the pore pressure, this quantity being termed the effective stress. If it is accepted that the mechanical behaviour of a material reflects the storage and dissipation of the power input, then it would be expected that this power input for a soil would depend on effective, not total stress. Such a result was obtained by Schofield and Wroth (1968) where they showed that, for an infinitely slow process, the mechanical work input to a soil is given by the product of the effective stress and the strain. (The result was obtained only for the special case of the triaxial test, but could be readily extended to more general stress states.)

Alternatively, if there is no deformation of the soil skeleton, the work input derives solely from the loss of excess pore pressure as the pore fluid seeps through the soil skeleton. The power input per unit volume may be calculated for this case as the product of the excess pore pressure gradient and the artificial seepage velocity.

The above two special cases represent extremes of soil behaviour: in the first, the entire power input is associated with the deformation of the soil skeleton, and in the second, the entire input is due to the viscous flow of the pore fluid. In many cases of engineering importance however, the two processes of deformation and seepage occur simultaneously, the best

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$$\int_V L dV = \int_V -(uw_j + \sigma_{ij}v_i)_{,j} + (\rho^{(w)}g_i w_i + \rho g_i v_i) dV$$

Again noting that since  $V$  is arbitrary this may be written in its local form, expansion of the differential then results in

$$L = -u_{,j}w_j - uw_{j,j} - \sigma_{ij,j}v_i - \sigma_{ij}v_{i,j} + \rho^{(w)}g_i w_i + \rho g_i v_i$$

Altering the dummy indices of the first term allows substitution of the definition of excess pore pressure gradient (4). Substitution of the compatibility condition (8) then allows rearrangement to

$$L = -u_{,i}'w_i - \sigma_{ij}v_{i,j} + uv_{j,j} - \sigma_{ij,j}v_i + \rho g_i v_i$$

The last two terms may be eliminated by use of the equilibrium condition (7). Using the relationship  $v_j = \delta_{ij}v_i$

$$L = -u_{,i}'w_i - \sigma_{ij}v_{i,j} + u\delta_{ij}v_{i,j}$$

Substituting the definition of effective stress (1)

$$L = -u_{,i}'w_i - \sigma_{ij}'v_{i,j}$$

Note that because of the symmetry of  $\sigma_{ij}'$ ,  $\sigma_{ij}'v_{[i,j]} = 0$ , so  $\sigma_{ij}'v_{i,j} = \sigma_{ij}'v_{(i,j)}$ . Making use of the symmetry of  $v_{(i,j)}$  to interchange indices, the substitution of the definition of strain rate (5) then yields the final result

$$L = \sigma_{ij}'\dot{\epsilon}_{ij} - u_{,i}'w_i$$

## DISCUSSION

It has been proven that there are no terms in the power input expression other than the two discussed above. The simple expression for the power input per unit volume therefore continues to apply for the case of finite deformation rate combined with seepage, the two terms being the product of the effective stress with the strain rate and the (negative) excess pore pressure gradient with the artificial seepage velocity. (The negative sign results simply from the sign convention for excess pore pressure gradient.)

This result may be used to give a new interpretation of the principle of effective stress. If the Terzaghi definition of effective stress is adopted it is observed that the total rate of work input per unit volume to the soil is given by the two terms  $(\sigma_{ij}'\dot{\epsilon}_{ij})$  and  $(-u_{,i}'w_i)$ . Clearly these may be interpreted as the rates of work input per unit volume to the soil skeleton and to the pore fluid respectively, and there is no coupling between the two processes of skeleton deformation and seepage. This result may be inverted to state that if there is no coupling between the work input to the soil skeleton and to the pore fluid, then the power input per unit volume to the soil skeleton is given at all times by the product of the effective stress with the strain rate. If it is further stated that the mechanical behaviour is simply a reflection of the way in which work is stored and dissipated, then if the processes of skeleton deformation and seepage are uncoupled, the mechanical behaviour of the skeleton will depend on the effective stress as defined by Terzaghi.

Although this gives an alternative interpretation of the principle of effective stress in terms of continuum mechanics, no statement is made here about whether soil would be expected to obey the principle, and hence show the uncoupling of the work terms. Any justification of the principle of effective stress for soils still rests on the arguments of particulate mechanics (notably Bishop, 1959) and the extensive body of experimental support for the theory. The

present analysis offers, however, a new interpretation in terms of continuum mechanics, in that the principle of effective stress is seen as a principle of the independence of the mechanical work input to the soil skeleton and to the pore fluid.

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